

Controlling Constraint Violation using Adjusted ADM/BSSN Systems

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Summary & Outlook

Formulation problem of Numerical Relativity

- Longer Evolutions by Adjusted ADM Systems
Teukolsky wave, 3+1 propagation
⇒ Standard ADM $\times 1.5 \sim 4$ life-time
- Trying to keep the Error at small value is better than to force the Error to zero.
- Longer Evolutions by Adjusted BSSN Systems
gauge wave / Gaudy model, 3+1 propagation
⇒ Standard BSSN $\times 10$ life-time
(See K.Kiuchi's poster)
- Next Step: Develop Auto-Control system of Lagrange multipliers

Formulation Problem?

Numerical Relativity

= Necessary for unveiling the strong nature of gravity

- GWs from NS-NS, BH-BH, NS-BH coalescence
- Relativistic phenomena like Cosmology, AGN, ...
- Mathematical feedbacks to Singularity, Exact Solutions, ...
- Laboratory of Gravitational theories: Higher dimensional models, ...

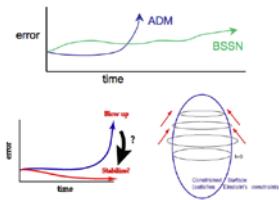
Current Standard Formulation

- BH-BH, NS-BH simulations (2005-...): Pretorius, UTB/Rochester, NASA, PSU, LSU, Jena, ...
- BSSN formulation: lapse function: 1+log slicing, shift vector: Gamma-freezing driver, initial data: puncture initial data

This combination works, anyway.

Why? Alternatives?

blow-up



For a review, please take a look
Shinkai & Yoneda, gr-qc/0209111

Adjusted Systems

General Procedure

- prepare a set of evolution eqs. $\partial_t u^a = f(u^a, \partial_b u^a, \dots)$
- add constraints in RHS $\partial_t u^a = f(u^a, \partial_b u^a, \dots) + F(C^a, \partial_b C^a, \dots)$
- choose appropriate $F(C^a, \partial_b C^a, \dots)$ to make the system stable evolution

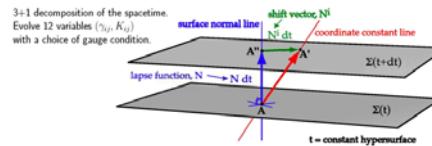
How to specify $F(C^a, \partial_b C^a, \dots)$?

- prepare constraint propagation eqs. $\partial_t C^a = g(C^a, \partial_b C^a, \dots)$
- and its adjusted version $\partial_t C^a = g(C^a, \partial_b C^a, \dots) + G(C^a, \partial_b C^a, \dots)$
- Fourier transform and evaluate eigenvalues $\partial_t C^a = A(C^a) C^a$

$t+0_+$ appropriate adjustments ⇒ Better Stability

The Standard ADM Formulation

(Arnowitt-Deser-Misner, 1962; York 1978)



	Maxwell eqs.	ADM Einstein eq.
constraints	$\text{div } E = -\epsilon p$ $\text{div } B = 0$	$\partial_t R + (\text{tr} K)^2 - K_{ij} K^{ij} - 2\kappa p_H + 2\lambda$ $D_i K_{ij}' - D_j K_{ii}' = \delta_{ij}$
evolution eqs.	$\frac{1}{c} \partial_t E = -\text{rot } B - \frac{4\pi}{c} J$ $\frac{1}{c} \partial_t B = -\text{rot } E$	$\partial_t K_{ij} = -2N K_{ij} + D_j N_i + D_i N_j$ $\partial_t K_{ij}' = N(\partial_t R_j + \text{tr} K K_{ij}) - 2N K_{ij} K_{ij}' - D_i D_j N + (D_i N^m) K_{mj} + (D_m N^i) K_{ij} - N^m D_m K_{ij} - N \gamma_{ij} \Lambda$ $- \kappa \alpha (S_{ij} + \frac{1}{2} \gamma_{ij} (p_H - \text{tr} S))$

BSSN Formulation

(Nakamura et al., 1987; Shibata-Nakamura 1995; Baumgarte-Shapiro 1996)

- define new variables $(\gamma_{ij}, K_{ij}, \lambda)$ instead of the ADM's $(\Sigma_{ij}, K_{ij}, \lambda)$, where $\lambda = \epsilon^{-1} K_{ij} - (1/2)\gamma_{ij} K_{ij}$, $\Gamma = P_{ij} K_{ij}$.
- use evolution constraint in Γ to keep, and impose $d\Gamma/dt = 1$ during the evolution.
- the set of evolution equations become:
 - $(\partial_t - L_{\Gamma}) \gamma_{ij} = -2K_{ij} + (1/2)\gamma_{ij} K_{ij}^2 - \gamma_{ij}^2 \nabla^2 \gamma_{ij}$
 - $(\partial_t - L_{\Gamma}) K_{ij} = -\epsilon^{-1} \nabla^2 \nabla_{ij} \gamma_{ij} - (1/2)\gamma_{ij} K_{ij}^2 + \epsilon^{-1} \alpha (\gamma_{ij} K_{ij})^2 + \epsilon^{-1} \alpha (\gamma_{ij} K_{ij})^2 - 2\kappa K_{ij} - 2\lambda K_{ij}^2$
 - $\partial_t \Gamma = -2(\beta \rho) \gamma_{ij}^2 - (1/2)\alpha (\gamma_{ij} K_{ij})^2 + (2\lambda K_{ij})^2 - 2\alpha K_{ij} \gamma_{ij}^2 - 2\alpha^2 \gamma_{ij}^2 \nabla^2 \gamma_{ij} - \partial_t (\beta \rho) \gamma_{ij}^2 - \gamma_{ij}^2 (\partial_t \alpha) \gamma_{ij}^2 + (2\lambda K_{ij})^2 \gamma_{ij}^2$

$$R_{ij} := \partial_t \gamma_{ij}^2 - \partial_t \gamma_{ij}^2 + \Gamma_{ij}^2 - \Gamma_{ij}^2 - \partial_t \gamma_{ij} \partial_t \gamma_{ij}$$

$$K_{ij}' := -2\beta \rho \partial_t \gamma_{ij} - 2\beta \rho \Gamma \partial_t \gamma_{ij} + \Gamma \partial_t \gamma_{ij} - \partial_t \beta \rho \partial_t \gamma_{ij}$$

$$\bar{R}_{ij} := -(1/2) \gamma_{ij}^2 \partial_t \gamma_{ij} + \beta \rho \partial_t \gamma_{ij}^2 + (1/2) \gamma_{ij}^2 \Gamma_{ij} \partial_t \gamma_{ij} + \beta \rho \Gamma_{ij}^2 \partial_t \gamma_{ij}$$

Guidelines for Better Formulation

Eigenvalue-analysis of Constraint Propagation eqs.

Conjecture on Constraint Amplification Factors (CAFs):

$$\partial_t \begin{pmatrix} \tilde{C}_1 \\ \vdots \\ \tilde{C}_N \end{pmatrix} = \begin{pmatrix} \text{Constraint Propagation Matrix} \\ \vdots \\ \tilde{C}_N \end{pmatrix} \begin{pmatrix} \tilde{C}_1 \\ \vdots \\ \tilde{C}_N \end{pmatrix}$$

We see more stable evolution, if CAFs have (A) negative real-part (the constraints are forced to be diminished), or (B) non-zero imaginary-part (the constraints are propagating away).

Adjusted ADM ADM vs BSSN Adjusted BSSN

Standard ADM has constraint violating mode!

- Better ADM must be available This Work
- Better BSSN must be available Kluchi's poster

Adjusted ADM formulation (1)

We adjust the standard ADM system using constraints as:

$$\partial_t \gamma_{ij} = -2\alpha K_{ij} + \nabla_i \beta_j + \nabla_j \beta_i + \nabla_i \beta_j + \nabla_j \beta_i + \nabla_i^2 \partial_t \gamma_{ij} + q^{ij} \partial_t (\nabla_i \partial_j \mathcal{H}), \quad (1)$$

$$\partial_t K_{ij} = \alpha R_{ij}^0 + \alpha K_{ij} - 2\alpha K_{ik} K_{kj}^i - \nabla_i \nabla_j \alpha + (\nabla_i \partial_j^0) K_{ij} + (\nabla_j \partial_i^0) K_{ij} + \beta^0 \nabla_k K_{ij}, \quad (2)$$

$$+ q_{ij} \mathcal{H} + S^k_{ij} M_k + r^k_{ij} (\nabla_k \mathcal{H}) + s^{kl}_{ij} (\nabla_k \partial_l \mathcal{M}_k), \quad (3)$$

with constraint equations

$$H := R^0 + K^2 - K_{ij} K^{ij}, \quad (5)$$

$$M_i := \nabla_j K_{ij} - \nabla_i K, \quad (6)$$

We can write the adjusted constraint propagation equations as

$$\partial_t \mathcal{H} = \{\text{original terms}\} + H_{\text{err}}^{(1)}[2] + H_{\text{err}}^{(2)}[\partial_t[2]] + H_{\text{err}}^{(3)}[\partial_t \partial_t[2]] + H_{\text{err}}^{(4)}[4], \quad (7)$$

$$\partial_t M_i = \{\text{original terms}\} + M_{\text{err}}^{(1)}[2] + M_{\text{err}}^{(2)}[\partial_t[2]] + M_{\text{err}}^{(3)}[\partial_t \partial_t[2]] + M_{\text{err}}^{(4)}[\partial_t[4]], \quad (8)$$

Adjusted ADM formulation (2)

Table 3: List of adjustments we made in the Schwarzschild spacetime. The entries are written in terms of (1) and (2). The entries "1" and "2" are the same as in table 1. The effects to amplitude factor, which is $\epsilon = 0$, are computed for each constraint system and the end frequency parts of 10%, respectively. The "N" term means that the numerical values are not available. The entries "no" and "yes" mean that the corresponding terms are not included and included, respectively, and thus "no" gives positive impacts, respectively. These differences are made in the $\epsilon = 0.001$ case. The values in parentheses are the values in the $\epsilon = 0.01$ case.

No.	Term	1	2	Adjustments	1	2	1	2	1	2	1	2
P-1	$\partial_t \gamma_{ij}$	no	no	no	no	no	no	no	no	no	no	no
P-2	$\partial_t K_{ij}$	no	no	no	no	no	no	no	no	no	no	no
P-3	$\partial_t \Gamma$	no	no	no	no	no	no	no	no	no	no	no
P-4	$\partial_t M_i$	no	no	no	no	no	no	no	no	no	no	no
P-5	$\partial_t \mathcal{H}$	no	no	no	no	no	no	no	no	no	no	no
P-6	$\partial_t M_k$	no	no	no	no	no	no	no	no	no	no	no
P-7	$\partial_t \mathcal{M}_k$	no	no	no	no	no	no	no	no	no	no	no
P-8	$\partial_t \mathcal{L}$	no	no	no	no	no	no	no	no	no	no	no
P-9	$\partial_t \mathcal{M}_i$	no	no	no	no	no	no	no	no	no	no	no
P-10	$\partial_t \mathcal{L}_k$	no	no	no	no	no	no	no	no	no	no	no
P-11	$\partial_t \mathcal{L}_i$	no	no	no	no	no	no	no	no	no	no	no
P-12	$\partial_t \mathcal{L}_k^2$	no	no	no	no	no	no	no	no	no	no	no
P-13	$\partial_t \mathcal{L}_i^2$	no	no	no	no	no	no	no	no	no	no	no
P-14	$\partial_t \mathcal{L}_k \mathcal{L}_i$	no	no	no	no	no	no	no	no	no	no	no
P-15	$\partial_t \mathcal{L}_k \mathcal{L}_k$	no	no	no	no	no	no	no	no	no	no	no
P-16	$\partial_t \mathcal{L}_i \mathcal{L}_i$	no	no	no	no	no	no	no	no	no	no	no
P-17	$\partial_t \mathcal{L}_k \mathcal{L}_k^2$	no	no	no	no	no	no	no	no	no	no	no
P-18	$\partial_t \mathcal{L}_i \mathcal{L}_i^2$	no	no	no	no	no	no	no	no	no	no	no
P-19	$\partial_t \mathcal{L}_k^2 \mathcal{L}_i$	no	no	no	no	no	no	no	no	no	no	no
P-20	$\partial_t \mathcal{L}_i^2 \mathcal{L}_k$	no	no	no	no	no	no	no	no	no	no	no
P-21	$\partial_t \mathcal{L}_k^2 \mathcal{L}_k$	no	no	no	no	no	no	no	no	no	no	no
P-22	$\partial_t \mathcal{L}_i^2 \mathcal{L}_i$	no	no	no	no	no	no	no	no	no	no	no
P-23	$\partial_t \mathcal{L}_k \mathcal{L}_i^2$	no	no	no	no	no	no	no	no	no	no	no
P-24	$\partial_t \mathcal{L}_i \mathcal{L}_k^2$	no	no	no	no	no	no	no	no	no	no	no
P-25	$\partial_t \mathcal{L}_k^2 \mathcal{L}_k^2$	no	no	no	no	no	no	no	no	no	no	no
P-26	$\partial_t \mathcal{L}_i^2 \mathcal{L}_i^2$	no	no	no	no	no	no	no	no	no	no	no
P-27	$\partial_t \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k$	no	no	no	no	no	no	no	no	no	no	no
P-28	$\partial_t \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i$	no	no	no	no	no	no	no	no	no	no	no
P-29	$\partial_t \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k$	no	no	no	no	no	no	no	no	no	no	no
P-30	$\partial_t \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i$	no	no	no	no	no	no	no	no	no	no	no
P-31	$\partial_t \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i$	no	no	no	no	no	no	no	no	no	no	no
P-32	$\partial_t \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k$	no	no	no	no	no	no	no	no	no	no	no
P-33	$\partial_t \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k$	no	no	no	no	no	no	no	no	no	no	no
P-34	$\partial_t \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i$	no	no	no	no	no	no	no	no	no	no	no
P-35	$\partial_t \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k$	no	no	no	no	no	no	no	no	no	no	no
P-36	$\partial_t \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i$	no	no	no	no	no	no	no	no	no	no	no
P-37	$\partial_t \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k$	no	no	no	no	no	no	no	no	no	no	no
P-38	$\partial_t \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i$	no	no	no	no	no	no	no	no	no	no	no
P-39	$\partial_t \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i$	no	no	no	no	no	no	no	no	no	no	no
P-40	$\partial_t \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k$	no	no	no	no	no	no	no	no	no	no	no
P-41	$\partial_t \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k$	no	no	no	no	no	no	no	no	no	no	no
P-42	$\partial_t \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i$	no	no	no	no	no	no	no	no	no	no	no
P-43	$\partial_t \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k$	no	no	no	no	no	no	no	no	no	no	no
P-44	$\partial_t \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i$	no	no	no	no	no	no	no	no	no	no	no
P-45	$\partial_t \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k$	no	no	no	no	no	no	no	no	no	no	no
P-46	$\partial_t \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i$	no	no	no	no	no	no	no	no	no	no	no
P-47	$\partial_t \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i$	no	no	no	no	no	no	no	no	no	no	no
P-48	$\partial_t \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k$	no	no	no	no	no	no	no	no	no	no	no
P-49	$\partial_t \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k \mathcal{L}_k$	no	no	no	no	no	no	no	no	no	no	no
P-50	$\partial_t \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i \mathcal{L}_i$	no	no	no	no	no	no	no	no	no	no	no
P-51	$\partial_t \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k$	no	no	no	no	no	no	no	no	no	no	no
P-52	$\partial_t \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i \mathcal{L}_k \mathcal{L}_i</math$											